The Crystal and Band Electronic Structures of Bis[4,5-ethylenedithio-4',5'-(vinylenedithio)tetrathiafulvalenium] Hexafluorophosphate, and Hexafluoroarsenate [EVT₂PF₆ and EVT₂AsF₆]

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The molecular and crystal structures of EVT_2PF_6 and EVT_2AsF_6 [EVT=4,5-ethylenedithio-4',5'-(vinylenedithio)tetrathiafulvalene] were determined by an X-ray diffraction method. Tight-binding band-structure calculations were made, and the results were compared with those of closely resembling conducting salts, α -(BEDT-TTF) $_2PF_6$ and VT_2PF_6 [BEDT-TTF=bis(ethylenedithio)tetrathiafulvalene; VT=bis(vinylenedithio)tetrathiafulvalene]. It was concluded that the extent of the two-dimensional nature of the band electronic structure increased systematically in the order VT_2PF_6 , EVT_2PF_6 , and α -(BEDT-TTF) $_2PF_6$; this result was analysed in terms of their crystal structures.

Studies concerning conducting organic salts have led to various organic superconductors during the last few years.1) The properties of conducting salts are closely related to their crystal structures. Although the existence of organic superconductors with twodimensional molecular arrangements has been established,2) most of the conducting salts have a stack of donor and/or acceptor molecules in the crystal to form a columnar structure. It has now been recognized that an intercolumnar interaction through chalcogenchalcogen contacts shorter than the van der Waals distance plays an important role in realizing the twodimensional electrical properties of conducting salts. This is believed to be effective for suppressing the Peierls transition inherent in a one-dimensional system. In order to determine the effect of a slight modification of a donor molecule on both the crystal and band electronic structures of a conducting salt, it is important to compare the crystal and band electronic structures of salts possessing similar molecular structures at the donor site. Knowledge obtained in this way will be useful in designing new conducting organic salts. For this purpose we noticed bis(ethylenedithio)tetrathiafulvalene (BEDT-TTF), 4,5-ethylenedithio-4',5'-(vinylenedithio)tetrathiafulvalene (EVT),3) and bis(vinylenedithio)tetrathiafulvalene (VT)4) to be donor molecules possessing similar molecular struc-We carried out X-ray crystallographic studies of EVT₂PF₆ and EVT₂AsF₆. Band electronic calculations were made based on a tight-binding method,

using the molecular orbitals obtained by an extended Hückel method.⁵⁾ These results were compared with those of α -(BEDT-TTF)₂PF₆ and VT₂PF₆, which have the same crystal system and space group (triclinic, PI)^{6,7)} as those of EVT₂PF₆ and EVT₂AsF₆. We found that a slight modification of the donor site results in a systematic change of both the crystal and band electronic structures.⁸⁾

Experimental

Electrochemical Crystallization. EVT was synthesized by a previously reported method.³⁾ Single crystals of EVT₂PF₆ and EVT₂AsF₆ were grown using an electrochemical crystallization method in tetrahydrofuran and chlorobenzene, respectively, under a galvanostatic condition ($\sim 1 \mu A$), using platinum wires as electrodes and tetrabutylammonium hexafluorophosphate or tetrabutylammonium hexafluoroarsenate as electrolytes.

Crystal Data. EVT₂PF₆: $(C_{10}H_6S_8)_2$ PF₆, M=910.33, triclinic, PI, a=6.47(2), b=7.80(3), c=15.81(5) Å, $\alpha=94.8(8)$, β =80.2(3), γ =100.2(6)°, V=772(5) ų, Z=1, D_c =1.96 g cm⁻³, $\mu(\text{Mo }K_{\alpha})=11.8 \text{ cm}^{-1}; \text{ EVT}_2\text{AsF}_6 (\text{C}_{10}\text{H}_6\text{S}_8)_2\text{AsF}_6, M=954.28,$ triclinic, PI, a=6.494(1), b=7.760(2), c=16.039(4) Å, $\alpha=$ 94.97(2), β =98.63(1), γ =79.58(2)°, V=784.5(3) ų, Z=1, D_c = 2.02 g cm⁻³, $\mu(\text{Mo }K_{\alpha})=21.6 \text{ cm}^{-1}$. The crystals were needle shaped with approximate dimensions of 0.78×0.21×0.08 mm³ for EVT₂PF₆ and 0.53×0.28×0.10 mm³ for EVT₂AsF₆. X-Ray diffraction data were collected by the θ -2 θ scan technique up to $2\theta=47^{\circ}$ and 50° for EVT₂PF₆ and EVT₂AsF₆ respectively on a Rigaku AFC-4 four-circle diffractometer, using graphite monochromatized Mo K_{α} radiation (λ = 0.71069 Å). The scan rates (θ) were 2° min⁻¹ for EVT₂PF₆ and 4° min⁻¹ for EVT₂AsF₆, and the scan widths $\Delta\theta$ = $(1.2+0.5 \tan \theta)^{\circ}$ for EVT₂PF₆ and $\Delta \theta = (1.0+0.5 \tan \theta)^{\circ}$ for EVT₂AsF₆. Backgrounds were counted for 5 s at both ends of each scan. No significant intensity decay of three standard reflections, which were measured after every 100 reflections, were detected for either crystal. For EVT₂PF₆, of the 2523 reflections measured, 1778 were observed ($|F_o| > 2\sigma(F_o)$). For EVT₂AsF₆, of the 2990 reflections measured, 2565 were observed ($|F_o| > 2\sigma(F_o)$). The data were corrected for both

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Lorentz and polarizaion factors, but not for absorption.

The structures were solved by a direct method (MULTAN 78),⁹⁾ and were refined anisotropically by a block-diagonal least-squares procedure (HBLS V) for nonhydrogen atoms.¹⁰⁾ Not all of the hydrogen atoms could be reasonably refined. Thus, they were relocated at the calculated positions with the isotropic temperature factors set equal to those of bonded carbon atoms. They were included in the calculations but were not refined. The minimized function

Table 1. Atomic Parameters of Nonhydrogen Atoms of EVT₂PF₆, with Equivalent Isotropic Temperature Factors¹²⁾

		iiperature raet		
Atom	x	y	z	$B_{ m eq}/{ m \AA}^2$
C(1)	0.492(2)	0.284(2)	0.5181(8)	2.9
C(2)	0.324(3)	0.151(2)	0.3864(9)	2.9
C(3)	0.532(3)	0.172(2)	0.3609(8)	3.0
C(4)	0.251(3)	-0.005(3)	0.233(1)	5.4
C(5)	0.453(3)	0.092(3)	0.1958(9)	4.2
C(6)	0.525(2)	0.338(1)	0.6009(8)	2.2
C(7)	0.476(2)	0.429(2)	0.7616(8)	2.8
C(8)	0.683(2)	0.461(2)	0.7356(7)	2.1
C(9)	0.537(3)	0.361(2)	0.9149(8)	3.4
C(10)	0.749(3)	0.399(2)	0.8942(8)	3.2
S(1)	0.2399(6)	0.2125(5)	0.4932(3)	3.4
S(2)	0.6982(5)	0.2619(5)	0.4374(2)	3.1
S(3)	0.1193(6)	0.0715(6)	0.3288(3)	3.9
S(4)	0.6658(6)	0.1268(6)	0.2597(3)	4.0
S(5)	0.3166(5)	0.3484(5)	0.6836(2)	3.0
S(6)	0.7766(6)	0.4044(5)	0.6299(2)	2.9
S(7)	0.3599(6)	0.4554(5)	0.8664(3)	3.2
S(8)	0.8696(6)	0.5257(5)	0.8068(3)	3.4
P	0.	0.	0.	2.8
$\mathbf{F}(1)$	0.014(3)	0.201(1)	0.0221(8)	7.8
F(2)	-0.231(1)	-0.012(2)	-0.0213(7)	6.3
F(3)	0.091(2)	0.042(1)	-0.0957(6)	6.0

Table 2. Atomic Parameters of Nonhydrogen Atoms of EVT_2AsF_6 , with Equivalent Isotropic Temperature Factors¹²⁾

Atom	x	у	z	$B_{ m eq}/ m \AA^2$
C(1)	0.5103(7)	0.2812(6)	0.5211(3)	3.0
C(2)	0.6771(7)	0.1498(6)	0.3899(3)	3.0
C(3)	0.4687(7)	0.1710(6)	0.3641(3)	3.1
C(4)	0.753(1)	-0.010(2)	0.2363(4)	6.4
C(5)	0.5489(8)	0.0923(7)	0.2004(3)	4.2
C(6)	0.4759(7)	0.3398(6)	0.6010(3)	2.8
C(7)	0.5221(7)	0.4333(6)	0.7591(3)	2.7
C(8)	0.3148(7)	0.4601(6)	0.7351(3)	2.8
C(9)	0.4610(8)	0.3674(7)	0.9131(3)	3.6
C(10)	0.2506(9)	0.4059(7)	0.8913(3)	4.1
S(1)	0.7610(2)	0.2141(2)	0.4940(1)	3.5
S(2)	0.3041(2)	0.2623(2)	0.4391(1)	3.4
S(3)	0.8839(2)	0.0699(2)	0.3313(1)	4.2
S(4)	0.3384(2)	0.1227(2)	0.2639(1)	4.3
S(5)	0.6831(2)	0.3508(2)	0.6821(1)	3.3
S(6)	0.2233(2)	0.4054(2)	0.6298(1)	3.1
S(7)	0.6390(2)	0.4593(2)	0.8640(1)	3.7
S(8)	0.1280(2)	0.5299(2)	.0.8049(1)	3.8
As	0.	0.	0.	3.2
$\mathbf{F}(1)$	-0.0076(9)	0.2184(5)	0.0208(3)	8.7
F(2)	0.2481(5)	-0.0220(7)	-0.0259(3)	7.5
F(3)	-0.1007(6)	0.0351(6)	-0.1027(2)	6.4

was $\Sigma \omega(\Delta F)^2$. The weighting functions used in the final stage of the refinements were $\omega = [\sigma^2(F_o) - 0.0045|F_o| + 0.0149|F_o|^2]^{-1}$ for EVT₂PF₆ and $\omega = [\sigma^2(F_o) - 0.0055|F_o| + 0.0029|F_o|^2]^{-1}$ for EVT₂AsF₆. The atomic scattering factors were taken from those of International Tables of X-Ray Crystallography. The final R indices were 0.118 for EVT₂PF₆ and 0.053 for EVT₂AsF₆. The atomic coordinates and equivalent isotropic temperature factors 2.13)

Electrical Conductivities. The temperature dependences of the electrical conductivities of the single-crystals of EVT_2PF_6 and EVT_2AsF_6 were measured by a four-probe method, using a D-type Cryomini Refrigerator (Osaka Sanso Co. Ltd).

Results and Discussion

Electrical Conductivities. Figure 1 shows the temperature dependence of the electrical conductivities of EVT₂PF₆ and EVT₂AsF₆ along the *b*-axis. The electrical conductivities of EVT₂PF₆ and EVT₂AsF₆ at room temperature were 0.76 S cm⁻¹ and 0.17 S cm⁻¹, respectively. They showed a semiconducting behavior with small activation energies of ca. 0.14 eV for EVT₂PF₆ and 0.11 eV for EVT₂AsF₆. α-(BEDT-TTF)₂PF₆ also showed a semiconducting behavior with a similar band gap of about 0.1 eV.⁶⁾ On the

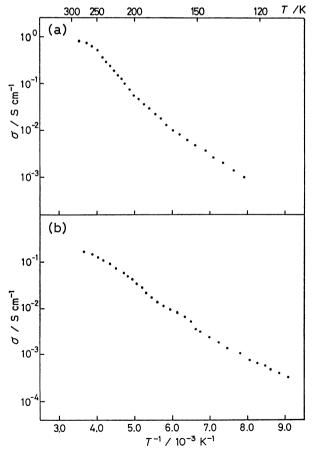
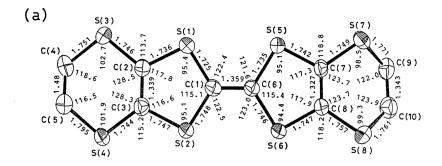


Fig. 1. Temperature dependences of the electrical conductivities of the single crystals of EVT₂PF₆ and EVT₂AsF₆ measured by a four probe method. (a) EVT₂PF₆, (b) EVT₂AsF₆.



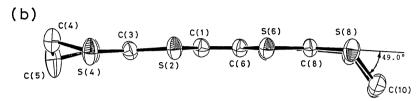


Fig. 2. The molecular structure of EVT molecule in EVT₂AsF₆, and the bond lengths and bond angles. (a) Over view, (b) side view. Nonhydrogen atoms are drawn as thermal ellipsoids with 50% probability level. Estimated standard deviations of the bond lengths and angles are 0.005—0.02 Å and 0.3—0.6° respectively.

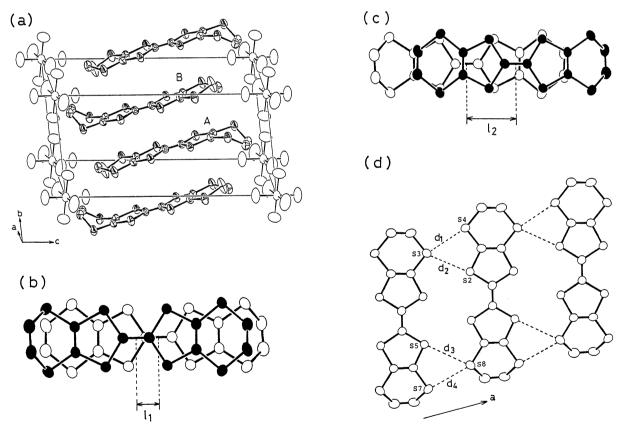
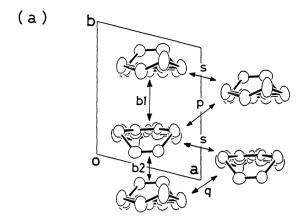
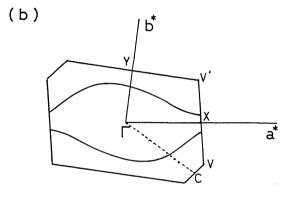


Fig. 3. (a) The crystal structure of EVT₂AsF₆. The molecules A and B are related by the symmetry operation (1-x, 1-y, 1-z). (b) The mode of a stacking of EVT molecules related by the symmetry operation (1-x, 1-y, 1-z). (c) The mode of a stacking of EVT molecules related by the symmetry operation (1-x, -y, 1-z). (d) The arrangement of the EVT molecules along the a-axis. The dotted lines indicate the S···S contacts shorter than the van der Waals distance $(3.7 \,\text{Å})$. The values of l_1 , l_2 , d_1 , d_2 , d_3 , and d_4 are summarized in Table 3.

other hand, VT₂PF₆ showed a metallic behavior above 180 K, below which the conductivity was activated $(E_a \sim 0.055 \text{ eV} (T < 100 \text{ K})).^{7)}$

Molecular and Crystal Structures of EVT₂AsF₆ and **EVT₂PF₆.** Figure 2 shows the molecular structure of the EVT molecule in the EVT₂AsF₆ crystal, together with the bond lengths and bond angles. The planarity of the EVT molecule was increased in its semication state (EVT^{+0.5}).3) No conformational disorder was found at the ethylene group of the EVT molecule. Figure 3 shows the crystal structure of EVT₂AsF₆. Only one EVT molecule is crystallographically independent, and the AsF₆⁻ anion is on the center of symmetry. The EVT molecules are stacked along the b-axis, as is usually observed for the conducting salts of TTF derivatives (Fig. 3(a)). Two kinds of the molecular overlap shown in Figs. 3(b) and 3(c) are observed. In one overlap mode (Fig. 3(b)), two EVT molecules are slightly shifted relative to each other within the unit cell. In the other overlap mode (Fig. 3(c)), the five-membered ring of one EVT molecule is overlapped with the six-membered ring of another EVT molecule across the unit cell. These two overlap modes are characterized by l_1 and l_2 in Figs. 3(b) and 3(c). No intermolecular S....S contact shorter than the van der Waals distance (3.7 Å) was found within the donor column. This fact is in contrast with the crystal structures of α-(BEDT-TTF)₂PF₆ and VT₂PF₆; intermolecular S····S contacts shorter than 3.7 Å were found within the donor column for these salts.^{6,7)} Table 3 shows the relative shifts of the donor molecules (l_1 and l_2) and intercolumnar S····S distances $(d_1, d_2, d_3, \text{ and } d_4) \text{ in } \alpha\text{-(BEDT-TTF)}_2\text{PF}_6, \text{ EVT}_2\text{PF}_6,$ EVT₂AsF₆, and VT₂PF₆ for the sake of a comparison (see also Figs. 3(b), 3(c), and 3(d)). All of the S····S contacts along the a-axis were shorter than 3.7 Å in EVT₂AsF₆, as shown in Fig. 3(d) and Table 3. Both α-(BEDT-TTF)₂PF₆ and VT₂PF₆ also possess intercolumnar S....S contacts shorter than 3.7 Å.6,7) The molecular and crystal structures of EVT2PF6 were found to be very similar to those of EVT₂AsF₆. Hence, the above arguments regarding the molecular and crystal structures of EVT₂AsF₆ are also applicable to the case of EVT₂PF₆. However, the final R value





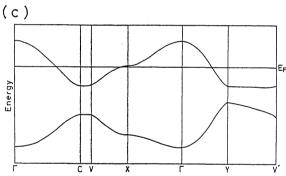


Fig. 4. (a) Intermolecular overlap integrals (b1, b2, p, q, and s), (b) Fermi surface, and (c) band structure of EVT₂PF₆. Since this salt shows a semiconducting behavior, the absence of the band gap (in other words, the appearance of the Fermi surface) will be caused by the neglect of on-site Coulomb repulsion among conducting electrons in the present calculation.

Table 3. Relative Shift of Donor Molecules (l_1 and l_2) and Intercolumnar S···S Distances between Donor Molecules (d_1 , d_2 , d_3 , and d_4) in the Crystals of α -(BEDT-TTF)₂PF₆, EVT₂PF₆, EVT₂AsF₆, and VT₂PF₆

Distances ^{a)} /Å	α-(BEDT-TTF) ₂ PF ₆	EVT ₂ PF ₆	EVT ₂ AsF ₆	$ m VT_2PF_6$
l_1	0.3	1.2	1.2	1.6
l_2	4.2	3.0	3.0	2.4
d_1	3.48	3.404(7)	3.394(3)	3.51
d_2	3.61	3.450(6)	3.503(2)	3.56
d_3	3.57	3.619(5)	3.664(2)	3.71
d_4	3.58	3.599(6)	3.592(2)	3.43

a) The distances l_1 , l_2 , d_1 , d_2 , d_3 , and d_4 are shown in Figs. 3(b), 3(c), and 3(d) for EVT₂AsF₆ as typical examples.

Table 4. Intermolecular Overlap Integrals of HOMO Calculated by an Extended Hückel Molecular Orbital Method^{a)}

Overlap integral	α -(BEDT-TTF) ₂ PF ₆	EVT ₂ PF ₆	EVT ₂ AsF ₆	VT_2PF_6
b1×10³	20.2	14.5	14.5	13.2
$b2 \times 10^{3}$	8.7	8.4	8.6	13.9
$p\times10^3$	2.9	1.8	1.6	2.1
$q \times 10^3$	0.6	3.0	3.3	1.3
s×10³	-2.5	-1.0	-0.8	-1.4

a) Overlap integrals b1, b2, p, q, and s are shown in Fig. 4 for EVT₂PF₆ as a typical example.

of the latter-mentioned salt was large (R=0.118) due to poor crystal quality.

Band Electronic Structures. Figure 4 shows intermolecular overlap integrals, Fermi surface, and the band structure of EVT₂PF₆ calculated by a tightbinding approximation based on an extended Hückel molecular orbital.¹⁴⁾ The band structure and Fermi surfaces of EVT₂AsF₆ were very similar to those of EVT₂PF₆.¹⁵⁾ We also obtained both closed and open Fermi surfaces for α -(BEDT-TTF)₂PF₆ and VT₂PF₆,⁷⁾ respectively. The appearance of the Fermi surfaces in α-(BEDT-TTF)₂PF₆, EVT₂PF₆, and EVT₂AsF₆ conflicts with their semiconducting properties. These contradictory results are caused by an on-site Coulomb repulsion among conducting electrons, which was not taken into account in the present calculation. Since the present calculation method gives a simple perspective concerning the dimensionality of many conducting organic salts,16) the following analysis is believed to be valid for any discussion about the relative band electronic structures of BEDT-TTF-, EVT-, and VT- salts. The values of intermolecular overlap integrals of HOMO levels in VT2PF6, EVT₂PF₆, EVT₂AsF₆ and α -(BEDT-TTF)₂PF₆ are summarized in Table 4. Here, the intermolecular overlap integrals (b1, b2, p, q, and s) in EVT₂PF₆ are shown in Fig. 4 as a typical example. The amount of the intermolecular transfer integral (t) is assumed to be of the order of $t \approx ES$ [E=orbital energy of HOMO] level (\sim -8.5 eV); S=intermolecular overlap integral]. Table 4 shows that a dimeric intermolecular interaction exists in the direction of the donor stack in α -(BEDT-TTF)₂PF₆; b1 is more than twice as large as b2. The ratio of b1/b2 decreases systematically in the order α -(BEDT-TTF)₂PF₆>EVT₂PF₆>VT₂PF₆. Thus, an almost homogeneous intermolecular interaction within the donor column was found for VT_2PF_6 (b1 \approx b2). This feature makes the effective band width of VT₂PF₆ the largest among these salts; this is related to the fact that only VT₂PF₆ shows a metal-like behavior. α -(BEDT-TTF)₂PF₆ much larger absolute value of the overlap integral |s| than those of the other salts; s is a measure of the sideby-side intermolecular interaction. The dimensionality of the band electronic structure of these salts is discussed below from the viewpoint of their conduction band widths.

Table 5. Relative Conduction Band Widths 17,18) of α -(BEDT-TTF) $_2$ PF $_6$, EVT $_2$ PF $_6$, EVT $_2$ AsF $_6$, and VT $_2$ PF $_6$

Salts	$\Delta W_{\Gamma ext{-}Y}$	$\Delta W_{\Gamma ext{-X}}$
α -(BEDT-TTF) ₂ PF ₆	18.6	17.0
EVT_2PF_6	22.8	13.6
EVT ₂ AsF ₆	23.8	13.0
VT_2PF_6	30.4	12.4

The matrix elements (H_{ij}) of a secular equation of α -(BEDT-TTF)₂PF₆, EVT₂PF₆, EVT₂AsF₆, and VT₂PF₆ are written as follows (see also Fig. 4):^{17,18)}

$$H_{11} = H_{22} = 2 t_s \cos(ka)$$

and

$$H_{12} = H_{21}^* = t_{b1} + t_{b2} e^{-ikb} + t_b e^{-ikc} + t_a e^{-i(kb+kc)}$$
.

Here, t_x (x=b1, b2, p, q, and s) denote the intermolecular transfer integrals. By solving the secular equation, we obtained the following equations for the conduction-band widths (ΔW), important measures for the dimensionality of the band electronic structures of the conducting salts:

$$\Delta W_{\Gamma-Y} = -2(t_{b2} + t_q)$$

and

$$\Delta W_{\Gamma-X} = 4 t_s - 2(t_p + t_q).$$

 $\Delta W_{\Gamma-Y}$ and $\Delta W_{\Gamma-X}$ are the conduction-band widths in the Γ -Y and Γ -X directions in the Brillouin zone (see Fig. 4).^{17,18)} Table 5 summarizes the relative values of the conduction band widths¹⁹⁾ ($\Delta W_{\Gamma-Y}$ and $\Delta W_{\Gamma-X}$) of α -(BEDT-TTF)₂PF₆, EVT₂PF₆, EVT₂AsF₆, and VT_2PF_6 . It is noticeable that $\Delta W_{\Gamma-Y}$ increases in the order α -(BEDT-TTF)₂PF₆<EVT₂PF₆<VT₂PF₆. On the other hand, $\Delta W_{\Gamma-X}$ changes in the reverse order. Since $\Delta W_{\Gamma-Y}$ and $\Delta W_{\Gamma-X}$ are comparable in α -(BEDT-TTF)₂PF₆, this salt is assumed to possess a two-dimensional electronic property. Table 5 shows that the extent of the two-dimensional nature is largest for α -(BEDT-TTF)₂PF₆, intermediate for EVT₂PF₆, and smallest for VT₂PF₆. The relation between the conduction band widths and the crystal structures are discussed below in more detail regarding these salts.

The $\Delta W_{\Gamma-X}$ value of α -(BEDT-TTF)₂PF₆ is the largest among those of the conducting salts given in Table

The main reason for this arises from the largest absolute overlap integral |s| of this salt (Table 4). All of the HOMO levels of BEDT-TTF, EVT, and VT have larger LCAO coefficients on the sulfur atoms than those on the carbon atoms. Thus, the overlap integral, s, is mainly determined by the side-by-side overlap between the lone-pair orbitals of the sulfur atoms belonging to the adjacent donor columns. The main reason that the $\Delta W_{\Gamma-Y}$ value for EVT₂PF₆ is larger than that for α-(BEDT-TTF)₂PF₆ arises from the larger value of q in the former salt. Here, qdenotes the intercolumnar molecular overlap integral between the donor molecules related by a symmetry operation (1-x, -y, -z). On the other hand, EVT salts have smaller values of $\Delta W_{\Gamma-X}$ than that of α -(BEDT-TTF)₂PF₆. Thus, the extent of the twodimensional nature of the former salts is less than that of the latter. In the case of VT_2PF_6 , $\Delta W_{\Gamma-Y}$ is much larger than those of the corresponding BEDT-TTF and EVT salts. The main reason for this arises from the largest value of b2 (intracolumnar molecular overlap integral). The mode of the the stacking of the donor molecules connected by the overlap integral b2 are characterized by l_2 in Table 3. In the case of α -(BEDT-TTF)₂PF₆, the mode of the intermolecular overlap is that of a "ring-external bond-type overlap" usually found in 1-D organic metals. On the other hand, a six-membered ring of one molecule overlaps with a five-membered ring of another molecule in EVT₂PF₆ and VT₂PF₆; the extent of the molecular overlap of the latter salts is larger than that of the former. These features are characterized by the order of l_2 values: α -(BEDT-TTF)₂PF₆>EVT₂PF₆>VT₂PF₆. Moreover, only VT₂PF₆ has four S····S contacts shorter than the van der Waals distance between the donor molecules connected by the overlap integral, b2. Since the terminal vinylene group of a VT molecule is less sterically crowded than the terminal ethylene group of a BEDT-TTF or an EVT molecule, especially in the presence of PF₆⁻ anion nearby, the VT molecules related by a symmetry operation (1-x, -y,-z) are stacked more closely than the latter molecules are in EVT₂PF₆ or α -(BEDT-TTF)₂PF₆; this gives rise to the largest value of b2 for the former salt. On the other hand, the intercolumnar overlap integrals (p, q,and s) of VT₂PF₆ are not very different from the corresponding values of the EVT salt; the $\Delta W_{\Gamma-X}$ value of VT₂PF₆ is similar to that of the EVT salt. Therefore, the extent of the one-dimensional band electronic property of VT₂PF₆ is the largest among the conducting salts listed in Table 5. Comparing the modes of the molecular overlaps characterized by l_1 values in Table 3, α-(BEDT-TTF)₂PF₆ showed a direct overlap of the molecules; i.e., l_1 is very small (~0.3 Å). On the other hand, EVT₂PF₆ and VT₂PF₆ showed appreciable shifts of the molecules. These molecular overlaps are related to the overlap integral, b1. Therefore, b1 of α -(BEDT-TTF)₂PF₆ is much larger than

those of EVT₂PF₆ and VT₂PF₆. However, b1 does not appear in the equations of the band widths, $\Delta W_{\Gamma-Y}$ and $\Delta W_{\Gamma-X}$, but does in the equations for the band gaps.

The electrical conductivities of various VT-salts have been measured, and most of them showed a transition.²⁰⁾ The metal-insulator temperature dependence of the X-ray crystallographic analysis of VT₂PF₆ revealed a doubling of the lattice constants below 170 K $[(a,b,c) \rightarrow (2a,2b,2c)]$. These results are assumed to be caused by the quasi-one-dimensional nature of the VT-salts. Since the extent of the twodimensional nature of the EVT-salts is expected to be greater than that of the VT-salts, as in the present study, and since EVT is the molecule with the slightest modification of BEDT-TTF,3) EVT-salts are expected to give conducting salts with electronic properties similar to those of BEDT-TTF salts. Syntheses of the other conducting EVT-salts are underway using electrochemical crystallization and a diffusion method.

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- 14) The positions of hydrogen atoms were estimated by assuming sp² and/or sp³ hybridizations of terminal carbons of a donor molecule for the computation of the extended Hückel molecular orbitals.
- 15) Since the band structures of EVT₂PF₆ and EVT₂AsF₆ are very similar, the band structure of EVT₂AsF₆ will not be discussed in detail in this paper.

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- 17) For α -(BEDT-TTF)₂PF₆ and VT₂PF₆, the expressions of the matrix elements can be used if kc is replaced by ka (see Refs 6 and 7).
- 18) For α -(BEDT-TTF)₂PF₆ and VT₂PF₆, $\Delta W_{\Gamma-X}$ should be interpreted as $\Delta W_{\Gamma-Z}$ (see Refs 6 and 7).
- 19) Since the transfer integrals t_x (x=b1, b2, p, q, and s) are of the order of ES (E=orbital energy of HOMO level) and the values of E are not very different from one another for BEDT-TTF, EVT, and VT, the relative conduction band widths are shown by adopting the overlap integrals multiplied by 10^3 in Table 4.
- 20) Unpublished data.